pression is $4 n^{2} /\left(\log _{2} n-\log _{2} \log _{2} n-1\right)$. The result stated at the beginning of this section now follows immediately.
4. Generalizations. The corresponding problem may be considered for $n$ by $n$ matrices whose entries are taken from the integers $0,1, \cdots, k-1$. If an operation on such a matrix consists of adding a multiple of some row to some other row modulo $k$, then it can be shown that the foregoing theorem remains valid in this more general situation for any fixed value of $k$. In fact, the bounds in Sections 2 and 3 will still hold if $\log _{2} n$ is replaced by $\log _{k} n$.

1. N. J. Fine \& I. Niven, "The probability that a determinant be congruent to a $(\bmod m), "$ Bull. Amer. Math. Soc., v. 50, 1944, pp. 89-93.

## Evaluation of $I_{n}(b)=2 \pi^{-1} \int_{0}^{\infty}\left(\frac{\sin x}{x}\right)^{n} \cos (b x) d x$ and of <br> Similar Integrals

## By Rory Thompson

Medhurst and Roberts [1] suggest the problem of evaluating $I_{n}(b)$ for nonintegral values of $b$. There will be developed in this note an effective recursion scheme for such a calculation. In particular, it can be used to evaluate $I_{n}(0)$ for moderate values of $n$.

Following a suggestion by Hamming [2, p. 164], we differentiate $I_{n}(b)$ with respect to the parameter $b$, which is permissible by virtue of uniform convergence of the resulting integral for $n>2$ and continuity of the corresponding integrand with respect to both $x$ and $b$.

Thus we obtain

$$
\begin{aligned}
I_{n}^{\prime}(b) & =-2 \pi^{-1} \int_{0}^{\infty}\left(\frac{\sin x}{x}\right)^{n-1} \sin x \sin (b x) d x \\
& =\pi^{-1} \int_{0}^{\infty}\left(\frac{\sin x}{x}\right)^{n-1}[\cos (b+1) x-\cos (b-1) x] d x \\
& =\frac{1}{2}\left[I_{n-1}(b+1)-I_{n-1}(b-1)\right] .
\end{aligned}
$$

If the first expression for $I_{n}{ }^{\prime}(b)$ is integrated by parts there results the relation

$$
\begin{aligned}
I_{n}^{\prime}(b)= & (n-1) b^{-1} 2 \pi^{-1} \int_{0}^{\infty}\left(\frac{\sin x}{x}\right)^{n} \cos (b x) d x \\
& -n b^{-1} 2 \pi^{-1} \int_{0}^{\infty}\left(\frac{\sin x}{x}\right)^{n-1} \cos x \cos (b x) d x \\
= & \frac{n-1}{b} I_{n}(b)-\frac{n}{2 b}\left[I_{n-1}(b+1)+I_{n-1}(b-1)\right] .
\end{aligned}
$$

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Elimination of $I_{n}{ }^{\prime}(b)$ between these expressions yields the recurrence relation

$$
I_{n}(b)=\frac{1}{2(n-1)}\left[(n+b) I_{n-1}(b+1)+(n-b) I_{n-1}(b-1)\right]
$$

This recursion scheme for calculating $I_{n}(b)$ is reasonably stable, owing to the addition of comparable positive numbers and to the fact that an error in either $I_{n-1}(b+1)$ or $I_{n-1}(b-1)$ is multiplied by at most $(n+b) / 2(n-1)$, which is less than unity, since $b<n$ for nonzero values of $I_{n}(b)$.

Starting values are readily given by the relations

$$
\begin{aligned}
I_{3}(b) & =\frac{1}{8}\left[(b+3)^{2}-3(b+1)^{2}\right], \quad 0 \leqq b \leqq 1 \\
& =\frac{1}{8}\left[(b+3)^{2}-3(b+1)^{2}+3(b-1)^{2}\right], \quad 1 \leqq b \leqq 3 \\
& =0, \quad b \leqq 3 .
\end{aligned}
$$

Furthermore, we observe that negative values of $b$ can be taken into account by use of the relation

$$
I_{n}(-b)=I_{n}(b)
$$

The iterative procedure just described was used to produce an 8 D table of $I_{n}(b)$ for $n=3(1) 100, b=0(0.1) 9$ in approximately 0.9 minute on an IBM 7094 system at the MIT Computation Center. This table has been deposited in the UMT file of this journal.

For integral values of $b$ the table was checked by using the recurrence formula in the form

$$
I_{n-1}(b+1)=\frac{2(n-1)}{n+b} I_{n}(b)-\frac{n-b}{n+b} I_{n-1}(b-1)
$$

This implies

$$
I_{n-1}(1)=\frac{n-1}{n} I_{n}(0) .
$$

Values of $I_{n}(0)$ may be obtained from 10D tables in [1] and [3], so that $I_{n}(b)$ can be evaluated for integral values of $b$ in a more compact form than the formulas in [1].

The use of recurrence formulas is applicable to the numerical evaluation of other integrals, including indefinite ones. One such example is the calculation of the chi-square distribution, which was accomplished by Harter [4] essentially by direct integration. In this case, integrating

$$
F_{n}(u)=\left[2^{n / 2} \Gamma\left(\frac{n}{2}\right)\right]^{-1} \int_{0}^{u} \exp \left(-\frac{x}{2}\right) x^{n / 2-1} d x
$$

by parts yields the recurrence relation

$$
F_{n}(u)=F_{n-2}(u)-H_{n}(u)
$$

where

$$
H_{n}(u)=u^{n / 2-1} /\left[2^{n / 2-1} \exp \left(\frac{u}{2}\right) \Gamma\left(\frac{n}{2}\right)\right]=\frac{u}{n-2} H_{n-1}(u)
$$

and

$$
\begin{aligned}
& H_{1}(u)=\exp \left(-\frac{u}{2}\right) \\
& F_{2}(u)=1-\exp \left(-\frac{u}{2}\right)
\end{aligned}
$$

This procedure for evaluating $F_{n}(u)$ is sufficiently fast to permit a direct search for percentage points, in lieu of interpolation. Thus eleven critical levels were calculated to $5 D$ for $n=2(2) 100$ in 1.8 minutes on an IBM 7094.

Many other types of integrals exist for which this recursion scheme is feasible, in particular, Fourier (and other) transforms similar to $I_{n}(b)$.

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## Evaluation of Some Integrals Involving the $\psi$-Function

By M. L. Glasser

In the Bateman manuscript project tables, Erdelyi et al. [1] list five integrals over the unit interval involving the $\psi$-function (logarithmic derivative of the gamma function). The first of these is trivial, the second is easily derived by integrating by parts to derive a differential equation in terms of the parameter $a$. The fourth and fifth formulas are obtained by equating the imaginary and real parts of the second and the third is simply the case $a=0$ of the fourth. The purpose of this note is to point out that this table can be easily extended by simple use of the properties of the $\psi$-function. For example, many convergent integrals of the form

$$
I=\int_{m}^{n} f(x) \psi(x) d x
$$

where $m$ and $n$ are integers and $f(x)$ is a function such that $f(x)=-f(m+n-x)$, can be evaluated exactly. Thus, by symmetry

$$
I=\frac{1}{2} \int_{m}^{n} f(x)\{\psi(x)-\psi(m+n-x)\} d x
$$

Now use of the relations $\psi(y+1)=\psi(y)+y^{-1}$ and $\psi(y)-\psi(1-y)=-\pi$ cot $\pi y$ gives immediately

